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## SPHERICAL INVOLUTES OF THE FIXED POLE CURVE $(C^*)$ ON THE TIMELIKE BERTRAND CURVE COUPLE

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Abstract: In this paper, it has been showed that every single indicatrix of tangents and indicatrix of binormals of the curve,  $\alpha^*$  are spherical involutes of the fixed pole curve,  $(C^*)$  by finding a transition link of the timelike Bertrand curve couple through Frenet frames.

Keywords: Lorentz Space, Timelike Bertrand Curve Couple, Spherical Involutes

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## **1.Preliminaries**

Let Minkowski 3-space  $IR_1^3$  be the vector space  $IR^3$  equipped with the Lorentzian inner product g given by

$$g(X,X) = x_1^2 + x_2^2 - x_3^2$$

where  $X = (x_1, x_2, x_3) \in IR^3$ . A vector  $X = (x_1, x_2, x_3) \in IR^3$  is said to be timelike if g(X, X) < 0, spacelike if g(X, X) > 0 and lightlike (or null) if g(X, X) = 0. Similarly, an arbitrary curve  $\alpha = \alpha(s)$  in  $IR_1^3$  where *s* is a arclength parameter, can locally be timelike, spacelike or null (lightlike), if all of its velocity vectors  $\alpha'(s)$  are respectively timelike, spacelike or null (lightlike) for every  $s \in IR$ . The norm of a vector  $X \in IR_1^3$  is defined by [2]

$$||X|| = \sqrt{|g(X,X)|}.$$

We denote by  $\{T(s), N(s), B(s)\}$  the moving Frenet frame along the curve  $\alpha$ . Let  $\alpha$  be a timelike curve with curvature  $\kappa$  and torsion  $\tau$ . Let frenet vector fields of  $\alpha$  be  $\{T, N, B\}$ .

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In this trihedron, T is timelike vector field, N and B are spacelike vector fields. Then Frenet formulas are given by [3]

$$T' = \kappa N \quad , N' = \kappa T - \tau B \quad , B' = \tau N \, . \tag{1}$$

Let  $\alpha$  be a timelike vector, the frenet vectors T timelike, N and B are spacelike vector, respectively, such that

$$T \times N = -B$$
,  $N \times B = T$ ,  $B \times T = -N$ 

and the frenet instantaneous rotation vector is given by [5]

$$W = \tau T - \kappa B$$
,  $||W|| = \sqrt{|\kappa^2 - \tau^2|}$ .

Let  $\varphi$  be the angle between W and -B vectors and if W is a spacelike vector, we can write

$$\begin{cases} \kappa = \|W\| \cosh \varphi \\ , \quad C = \sinh \varphi T - \cosh \varphi B \\ \tau = \|W\| \sinh \varphi \end{cases}$$
(2)

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Let  $X = (x_1, x_2, x_3)$  and  $Y = (y_1, y_2, y_3)$  be the vectors in  $IR_1^3$ . The cross product of X and Y is defined by [1]

$$X \wedge Y = (x_3y_2 - x_2y_3, x_1y_3 - x_3y_1, x_1y_2 - x_2y_1).$$

The curvatures drawn by unit speed non-null curve,  $\alpha: I \to IR_1^3$  at the point  $\alpha(s)$  with the frenet vectors T, N, B and the unit Darboux vector C over the Lorentzian unit sphere  $S_1^2$  or Hyperbolic unit sphere  $H_0^2$  are named respectively as indicatrix of tangents, indicatrix of principal normals , indicatrix of binormals and fixed pole curve. These curvatures are indicated in order as (T), (N), (B) and (C) [4].

Let  $\alpha: I \to IR_1^3$  and  $\alpha^*: I \to IR_1^3$  be two timelike curves and if the tangent of  $\alpha$ ,  $\alpha(s)$  posses through the point  $\alpha^*(s)$  and  $\langle T^*(s), T(s) \rangle = 0$ , then the curve  $\alpha^*$  is said to be the involute of  $\alpha$  [6].

2. Spherical Involutes of the Fixed Pole Curve  $(C^*)$  on the Timelike Bertrand Curve Couple

**Definition 2.1:** Let  $\{T, N, B\}$  and  $\{T^*, N^*, B^*\}$  be respectively the frenet frames of the timelike curves,  $\alpha: I \to IL^3$  and  $\alpha^*: I \to IL^3$  at points  $\alpha(s)$  and  $\alpha^*(s)$ . If the principal normal vectors N and  $N^*$  are linearly dependent, then the pair  $(\alpha, \alpha^*)$  is said to be timelike Bertrand curve couple.

**Theorem 2.1:** There is a connection between timelike Bertrand curve couple and Frenet frames that are written as followings

$$\begin{cases} T^* = -\cosh\theta T + \sinh\theta B\\\\N^* = N\\\\B^* = -\sinh\theta T + \cosh\theta B. \end{cases}$$

Here, the angle  $\theta$  is the angle between T and  $T^*$ .

**Proof:** By taking the derivative of  $\alpha^*(s) = \alpha(s) + \lambda N(s)$  with respect to arc lenght *s* and using the equation, we get

$$T^* \frac{ds^*}{ds} = T(1 + \lambda \kappa) - \lambda \tau B.$$
(4)

The inner products of the above equation with respect to T and B are respectively defined as

$$\begin{cases} -\cosh\theta \frac{ds^*}{ds} = 1 + \lambda\kappa \\ -\sinh\theta \frac{ds^*}{ds} = \lambda\tau \end{cases}$$

and substituting these present equations in (4) we obtain

$$T^* = -\cosh\theta T + \sinh\theta B.$$
<sup>(5)</sup>

Here, by finding the derivative of (4) and using (1) we get

$$N^* = N$$
.

Firstly, we can write

$$B^* = -\sinh\theta T + \cosh\theta B \tag{6}$$

by availing the equation  $B^* = -(T^* \times N^*)$ .

By the derivative of  $\alpha_{T^*}(s_{T^*}) = T^*(s)$  with respect to arc-lenght  $s_{T^*}$  parameter, we get

$$T_{T^*} = \frac{dT^*}{ds} \cdot \frac{ds}{ds_{T^*}}$$

Afterwards, by some algebraic manipulations and substituting (5) in  $T_{T^*}$ , the following result can be achieved

$$T_{T^*} = \mp N . \tag{7}$$

Similarly, by taking the derivative of  $\alpha_{B^*}(s_{B^*}) = B^*(s)$  with respect to arc-lenght  $s_{B^*}$  parameter, we get

$$T_{B^*} = \frac{dB^*}{ds} \cdot \frac{ds}{ds_{B^*}}.$$

By using the equation (6), we write down

$$T_{B^*} = \mp N \,. \tag{8}$$

Lastly, by taking the derivative of  $\alpha_{C^*}(s_{C^*}) = C^*(s)$  with respect to arc-lenght  $s_{C^*}$  parameter, we obtain

$$T_{C^*} = \frac{dC^*}{ds} \cdot \frac{ds}{ds_{C^*}}$$

If  $W^*$  is spacelike, then by considering (2) and with some algebraic operation we get

$$T_{c^*} = \cosh \varphi^* T^* - \sinh \varphi^* B^*, \tag{9}$$

if  $W^*$  is timelike, then then by considering (3) and with some algebraic operation we get

$$T_{c^*} = \sinh \varphi^* T^* - \cosh \varphi^* B^*. \tag{10}$$

**Theorem 2.2:** Let  $(\alpha, \alpha^*)$  be timelike Bertrand curve couple. Tach of the indicatrix of tangents  $(T^*)$  and the indicatrix of binormals  $(B^*)$  of  $\alpha^*$  curve is a spherical involute of the fixed pole curve  $(C^*)$ .

**Proof:** In order to show that  $(T^*)$  and  $(B^*)$  of the  $\alpha^*$  curve is each a spherical involute of  $(C^*)$ , we need to prove the following statements

$$\langle \frac{dT^*}{ds}, \frac{dC^*}{ds} \rangle = 0$$

and

$$\langle \frac{dB^*}{ds}, \frac{dC^*}{ds} \rangle = 0.$$

If  $W^*$  is spacelike, by taking into account (7) and (9) we write down

$$\langle \frac{dT^*}{ds}, \frac{dC^*}{ds} \rangle = \langle N, \cosh \varphi^* T^* - \sinh \varphi^* B^* \rangle.$$

Next, by using (5) and (6) and doing required manipulations, the following result can be obtained

$$\langle \frac{dT^*}{ds}, \frac{dC^*}{ds} \rangle = 0$$

Furthermore, by exploiting the equations (5), (6), (8) and (9), we do the similar calculations to get

$$\langle \frac{dB^*}{ds}, \frac{dC^*}{ds} \rangle = 0$$

If  $W^*$  is timelike, we use (7) and (10) to reach

$$\langle \frac{dT^*}{ds}, \frac{dC^*}{ds} \rangle = \langle N, \sinh \varphi^* T^* - \cosh \varphi^* B^* \rangle.$$

Here, by substitution of (5) and (6) in the above the formula, we write

$$\langle \frac{dT^*}{ds}, \frac{dC^*}{ds} \rangle = 0$$

Once again, when the given relations (5), (6), (8) and (10) are taken into consideration, we get

$$\langle \frac{dB^*}{ds}, \frac{dC^*}{ds} \rangle = 0$$

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